QCD at LHC: theoretical developments

(mostly within the last 1 year)

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Particle physics is driven by the belief that:





... are driven and described by the same microscopic forces

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Here is the particle physicist's picture of the world:



It is all about the desert; what is it – what's its nature?

Is it merely a desert?

Or an oasis?

Or perhaps a jungle?







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There are several important problems that are in the realm of particle physics:

Ex: Confinement:

- ✓ An outstanding problem.
- Yet we know how to go around it and keep making progress.

Proof: The LHC





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The Dark Matter Problem



Especially after WMAP it became clear that:





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Why did I bring Dark Matter into this discussion?

Dark matter is a different story:

We do not know how to solve it
And we do not know how to circumvent it ...

✓ It has to have some microscopic explanation

 (more subtle) If there is a jungle of particles in the desert, then such new physics offers Dark Matter candidates.

In a way, conceptually, New Physics implies a resolution to the dark matter problem.

The opposite is not quite true:

We should view the absence of bSM physics at the LHC, if it comes to that, as a strong guide for understanding the mystery of Dark Matter

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The modern physics at particle accelerators



We have had great successes at accelerator-based physics in the recent past

Discovered Higgs boson:





... established the CKM paradigm:



40 years of tireless scrutiny: no deviation from the SM so far

- The apparent success of the SM can hardly by overstated.
- Yet, there is much more to do!

The role of QCD at Colliders (LHC)

- QCD is everywhere (we collide hadrons, we measure hadrons; strong coupling is largest in SM)
- How QCD helps:
 - Increases the accuracy of SM predictions for signals that we care about
 - Higgs
 - PDF's and α_s
 - Vector bosons
 - Jets
 - Top quarks (incl. results on top quark mass)
 - QCD and BSM searches
- For the QCD aficionados (and, yes, there are many of them)
 - LO is long dead, live NLO! (despite the fact that I'll focus on NNLO)
 - Parton showers, resummation and all that

Lifting the accuracy of SM predictions



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Here is the big picture



• LHC Run I: impressively broad agreement with SM!

Higgs

• LHC Run I discovered the Higgs and established it is SM-like



Higgs

- We want to know as much as possible about the Higgs. This means precise SM predictions to compare with experiment.
- Most pressing question: the uncertainty of the total cross-section
- It necessitated the calculation of the N³LO correction (a first for hadron colliders!)

Anastasiou, Dulat, Duhr, Furlan, Gehrmann, Herzog, Lazopoulos, Mistlberger '15



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• Total cross-section at N³LO:

Claude Duhr, Zurich Workshop 2016

$\sigma[\mathrm{pb}]$	$\delta_{ m PDF}$	δ_{lpha_s}	$\delta_{ m scale}$	$\delta_{ m trunc}$	$\delta_{ ext{PDF-TH}}$	$\delta_{ m EW}$	δ_{tb}	δ_{1/m_t}
48.48	±0.90pb	±1.26pb	$^{+0.09}_{-1.11}\mathrm{pb}$	±0.12	±0.56	±0.48	±0.34	±0.48
	±1.86%	±2.60%	$^{+0.2}_{-2.3}\%$	±0.25%	±1.15%	±1.00%	±0.70%	±1.00%

- Uses NNLO pdf; no N³LO pdf's available (likely 1% effect) See also Forte et al '14
- EW corrections exact at NLO; at mixed QCD-EW included in an EFT approach (gauge bosons integrated out into Wilson coefficients)
- Quark masses (m_t m_b) included exactly at NLO. NNLO desirable
- Threshold resummation likely not pressing issue anymore.
- Basically, at N3LO the Higgs cross-sections starts to look just like the NNLO cross-sections of 2-to-2 processes (top-pair, for example)

Higgs couplings

ZPW 2016

G. Petrucciani (CERN)



Uncertainty breakdown



Theory errors are subdominant at present but in some cases are close contenders

Strong coupling α_s



- What happened?
 - Jump in the error from lattice.
 - Inclusion of ttbar measurement which is in downward fluctuation.

Strong coupling α_{s} α_{s} results from hadron collider data all hadron collider (except tbar) and HERA results in NLO LEP, PETRA and tbar in NNLO 0.25 α_{s} CMS R_{32} ratio \rightarrow HERA



- LHC data provides good access to α_{s} , albeit with larger error.
- Allows unprecedented access to running of α_s at high scales (TeV) from, for example, jets and ttbar.

Parton distribution functions

• New generation of global PDF sets available: CT14, NNPDF3.0 and MMHT14.



- They are compatible with the PDF4LHC recommendation and their fluxes are rather similar.
- Some other sets differ (see above).
- Essential improvements are expected once LHC top differential calculations (NNLO now available) as well as jet calculations (NNLO to appear soon) are included.

Parton distribution functions

New combination set: PDF4LHC15:

See arXiv:1510.03865

Provides both MC and Hessian sets with varying number of members



Figure 8: Comparison of the MC900 PDFs with the sets that enter the combination: CT14, MMHT14 and NNPDF3.0 at NNLO. We show the gluon and the up, anti-down and strange quarks at Q = 100 GeV. Results are normalized to the central value of MC900.

Parton distribution functions

- Sets with EW corrections available (NNPDF2.3): makes possible the consistent calculation of mixed EW-QCD corrections up to NNLO in QCD.
- Soft-gluon resummation's effect on pdf studied (NNPDF3.0) in NNLO with DIS, DY and top data

See arXiv:arXiv:1507.01006



 For the SM Higgs, what matters is the resummation in the Higgs partonic cross-section; resummation in PDF is insignificant.

From Higgs, to Higgs + jets, to Higgs decays (VV+jets)



Higgs + 1 jet at NNLO

• Studied extensively by several groups (large m_t limit)

Boughezal, Caola, Melnikov, Petriello, Schulze '13 Chen, Gehrmann, Jaquier, Glover '14 Boughezal, Focke, Giele, Liu, Petriello '15 Caola, Melnikov, Schulze '15

- Moreover the Higgs boson has been decayed:
 - $H+j \rightarrow \gamma\gamma+j$
 - $H+j \rightarrow WW + j \rightarrow e \mu \nu \nu + j$
- Such calculations allow for precise predictions that directly match the experimental setup (and are thus very useful)
 See arXiv:1508.02684



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Higgs + up to 3 jets at NLO_{Cullen et al (GOSAM) 1307.4737}

• Very significant NLO corrections. Great reduction in theoretical uncertainty







From 1506.01016

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VV (vector boson pair) production at NNLO

- Vector boson pair production is motivated:
 - an alternative to/decay of Higgs boson
 - in its own right
- Thus all precision requirements for Higgs production are directly translated into VV.
- Tremendous progress has been achieved so far: since VV is a colorless final state, one can compute NNLO QCD corrections with effectively NLO methods Catani, Grazzini '07
- By now all relevant combinations of pairs of W,Z and γ are known to NNLO
- These were some of the early 2-to-2 NNLO calculations and showed surprisingly large NNLO corrections that were essential for finding agreement with data! (especially γγ).



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ZZ production with Z decay

• Improved $\Delta \phi$ distribution

Kallweit, Rathlev, Grazzini '15

• Data cannot discriminate the lepton PT yet but the NNLO K-factor has significant shape.



ZZ production with Z decay

• Fiducial cross-section:

Kallweit, Rathlev, Grazzini '15

Channel	$\sigma_{\rm LO}~({\rm fb})$	$\sigma_{\rm NLO}~({\rm fb})$	σ _{NNLO} (fb)	σ_{\exp} (fb)
$e^{+}e^{-}e^{+}e^{-}$	3.547(1)+2.9%	$5.047(1)^{+2.8\%}_{-2.3\%}$	5 79 (2)+3.4%	$4.6^{+0.8}_{-0.7}(\text{stat})^{+0.4}_{-0.4}(\text{syst.})^{+0.1}_{-0.1}(\text{lumi.})$
$\mu^+\mu^-\mu^+\mu^-$	3.347(1) - 3.9%		(2) - 2.6%	$5.0^{+0.6}_{-0.5}(\text{stat})^{+0.2}_{-0.2}(\text{syst.})^{+0.2}_{-0.2}(\text{lumi.})$
$e^+e^-\mu^+\mu^-$	$6.950(1)^{+2.9\%}_{-3.9\%}$	$9.864(2)^{+2.8\%}_{-2.3\%}$	$11.31(2)^{+3.2\%}_{-2.5\%}$	$11.1^{+1.0}_{-0.9}(\text{stat})^{+0.5}_{-0.5}(\text{syst.})^{+0.3}_{-0.3}(\text{lumi.})$

8 TeV

- Curiously, the NNLO corrections significantly improve agreement with ATLAS but the same flavor channels seems to go away from data ...
- However experimental errors are large but this will change at 13 TeV

Channel	$\sigma_{\rm LO}~({\rm fb})$	$\sigma_{\rm NLO}~({\rm fb})$	$\sigma_{\rm NNLO}~({\rm fb})$	σ_{\exp} (fb)	
$e^+e^-e^+e^-$	5 007(1)+4%	6.157(1)+2%	$7 14(2)^{+2\%}$	$8.4^{+2.4}_{-2.0}(\text{stat})^{+0.4}_{-0.2}(\text{syst.})^{+0.5}_{-0.3}(\text{lumi.})$	1
$\mu^+\mu^-\mu^+\mu^-$	$5.007(1)_{-5\%}$	$0.137(1)_{-2\%}$	7.14(2)-2%	$6.8^{+1.8}_{-1.5}(\text{stat})^{+0.3}_{-0.3}(\text{syst.})^{+0.4}_{-0.3}(\text{lumi.})$	
$e^+e^-\mu^+\mu^-$	$9.906(1)^{+4\%}_{-5\%}$	$12.171(2)^{+2\%}_{-2\%}$	$14.19(2)^{+2\%}_{-2\%}$	$14.7^{+2.9}_{-2.5}(\text{stat})^{+0.6}_{-0.4}(\text{syst.})^{+0.9}_{-0.6}(\text{lumi.})$	

WW production at NNLO

- Essential for understanding EWSB physics
- NNLO correction reduces tension with ATLAS; agrees with CMS

•



NNLO correction similar in size to $H \rightarrow WW^*$

Gehrmann, Grazzini, Kallweit et al '14

T	\sqrt{s} leV	σ_{LO}	σ_{NLO}	σ_{NNLO}	$\sigma_{gg \to H \to WW^*}$
	7	$29.52^{+1.6\%}_{-2.5\%}$	$45.16^{+3.7\%}_{-2.9\%}$	$49.04^{+2.1\%}_{-1.8\%}$	$3.25^{+7.1\%}_{-7.8\%}$
	8	$35.50^{+2.4\%}_{-3.5\%}$	$54.77^{+3.7\%}_{-2.9\%}$	$59.84^{+2.2\%}_{-1.9\%}$	$4.14^{+7.2\%}_{-7.8\%}$
-	13	$67.16^{+5.5\%}_{-6.7\%}$	$106.0^{+4.1\%}_{-3.2\%}$	$118.7^{+2.5\%}_{-2.2\%}$	$9.44^{+7.4\%}_{-7.9\%}$
-	14	$73.74^{+5.9\%}_{-7.2\%}$	$116.7^{+4.1\%}_{-3.3\%}$	$131.3^{+2.6\%}_{-2.2\%}$	$10.64^{+7.5\%}_{-8.0\%}$

- Hard to separate WW from top-pair production;
- b-jets essential in this:



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V+jet production at NNLO

- Vector boson production is the absolute classics (think Drell-Yan) at hadron collider
- Served as the discovery mode for W and Z bosons in 1983 at SPS
- First NNLO corrections to Drell-Yan were computed 25 years ago

Hamberg, van Neerven, Matsuura '91 Harlander, Kilgore '02

And differential vector boson production around 10 years ago

Anastasiou, Dixon, Melnikov, Petriello '03

- V+jet calculation needed in order to have full NNLO accuracy for the V P_T spectrum
- Recall: precision requirements in vector boson production is very high: it is at the-% level and was even proposed as a luminosity monitor for LHC.

Z+jet production at NNLO

Gehrmann-De Ridder, Gehrmann, Glover, Huss, Morgan '15 Boughezal, Campbell, Ellis, Focke, Giele, Liu, Petriello '15



- As expected, impressive impact on the PT distribution of Z-boson
- Notice the very high perturbative stability of the fiducial cross-section:

$$\sigma_{LO} = 103.6^{+7.7}_{-7.5} \text{ pb},$$

$$\sigma_{NLO} = 144.4^{+9.0}_{-7.2} \text{ pb},$$

$$\sigma_{NNLO} = 140.3^{+0.0}_{-1.4} \text{ pb}$$

W+jet production at NNLO

Boughezal, Focke, Liu, Petriello '15





Impact of NNLO corrections is significant, just as for Z+jet:

$p_T^{jet} > 30 \text{ GeV}, \eta_{jet} < 2.4$				
Leading order:	533^{+39}_{-38} pb			
Next-to-leading order:	$797^{+63}_{-49}~\rm{pb}$			
Next-to-next-to-leading order:	$787^{+0}_{-8} {\rm \ pb}$			

Fiducial cross-section at 1% level !

Pure QCD beasts: dijets and top-pair production



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Dijet production towards NNLO

- If there is one thing at hadron colliders that's a lot of jets!
- Measured over a large energy range and over many orders of magnitude
- Overall NLO QCD (+EW) agrees with data.





(Jets are essential in PDF's)

- What we really need is high-precision comparison, in order to:
 - Search for BSM physics decaying to jets
 - Provide input to SM: measurement of α_s and extraction of PDF's (including or not jet data in pdf's has been one of the most debated subjects in pdf community)

Dijet production towards NNLO

 The expectation is that NNLO QCD (supplemented with EW corrections) will be able to Dittmaier, Huss, Speckner '12

significantly increase the theoretical precision.

Expectation is based on several partial NNLO contributions (gg-> gg, qqbar-> gg)

A number of papers up to 2014: Currie, Gehrmann-De Ridder, Gehrmann, Glover, Pires

+ Ongoing Work





Looking forward to the complete result!





- Top discovered at the Tevatron but statistics there was very limited (~1k events)
- LHC gets the chance to produce lots of top events (>100k events recorded at Run I)
- LHC Run 2 cross-section larger by a factor of 4.
- The LHC should, for the first time, study the top completely, all its couplings and parameters.
- Top is (most) important background for most BSM searches.
- Interesting anomalies (top forward-backward asymmetry at the Tevatron)
- Important for SM Higgs
- So far the only NNLO input for gluon pdf from hadron colliders
- Measurement of α_S . Top mass is a major input when extending SM towards GUT scales

(think vacuum stability, Higgs inflation).

Impressive agreement for the total cross-section (level of 4-5%)



- Notable: after a month of data taking the largest error, by far, is the one due to luminosity!
- ✓ Cancels in the tt/Z ratio. Excellent agreement with NNLO SM.



• A lot of recent activity:

Czakon, Heymes, Fiedler, Mitov '15

- fully differential NNLO QCD production for stable top quarks at the Tevatron and LHC
- This can easily be combined with EW corrections (will be important for TeV scales)
- No top decay implemented at NNLO. Understood in principle. This is for the future.





Differential distributions at the LHC: important in the context of the "top P_T discrepancy"



• Several qualifications:

- Lepton- and jet-based observable appear to be fine.
- Top quark-level ones no so much.
- But tops are not measured; they are "inferred" from data using MC's.
- Therefore, any discrepancy between SM top quark predictions and 'measurements' are testing how well current MC's describe top production.
- Implications are very broad and go much beyond top physics: Higgs, BSM.

Top-pair production at NNLO: P_T spectrum

- NNLO QCD corrections systematically improve the agreement with CMS data.
- Agreement with ATLAS (not shown) even better.
- NNLO does what one normally expects:
 - Convergence
 - Decrease of scale error
 - Pdf error not included



Approximate results within the antenna approach have also appeared

Abelof, Gehrmann-De Ridder, Majer `15

Top-pair production at NNLO: M_{tt} spectrum



- ✓ The quality of the calculation is high:
 - ✓ Fine binning
- NNLO does what one normally expects:
 - Convergence
 - Decrease of scale error
 - Pdf error not included
 - Threshold effects can be seen
 - Note the extreme stability of the shape: no change from NLO to NNLO (0.5% or so)
 - An opportunity for searches?

Top-pair production at NNLO: PDF dependence



Figure 11. As in fig. 10 but for the normalised to unity distributions.

Normalized distributions show very small sensitivity to PDF's

 \checkmark Good news for m_{top} extractions from differential distributions.

Top quark mass

Not only is the top mass a fundamental SM parameter. It plays outsize role in extending the SM from current collider energies to GUT energies:



SM vacuum stability:

De Simone, Hertzbergy, Wilczek 0812.4946

Degrassi, Di Vita, Elias-Miro, Espinosa, Giudice, Isidori, Strumia '12

- The implications to BSM physics are well known.
- The notable fact is that 1 GeV change in the top mass completely alters the predictions.
- So, how well do we know the top mass anyway?

Top quark mass

- Look at the spread across current measurements:
 - > m_t = 173.34 ± 0.76 GeV [World Average]
 - $> m_t = 172.04 \pm 0.77 \text{ GeV} [CMS Collaboration]$
 - > m_t = 174.98 ± 0.76 GeV [D0 Collaboration]
- Comparable uncertainties; rather different central values!
 - Spread likely due to different theory systematics! Many methods proposed (recent reviews) Juste et al arXiv:1310.0799 Moch et al arXiv:1405.4781
- I would single out leptonic observables since they are cleaner and, supposedly, under better theory control
 Kawabata, Shimizu, Sumino, Yokoya '11-`14 Frixione, Mitoy `14
- Another important issue: how well can we determine the top mass at LHC?
 - Current m_t error of O(1GeV) could in principle go down even below O(100MeV). Therefore, pole mass calculations for the LHC are fine.
 - Finally, what is the ultimate precision on m_t one might expect?
 - 50-100MeV from a threshold scan at a linear e⁺e⁻ collider.
 - N³LO corrections recently completed

Beneke, Kiyo, Marquard, Penin, Piclum, Steinhauser 15

QCD and BSM



Few thoughts on BSM and QCD

- Not only is QCD important for BSM, but BSM plays great role in developing QCD applications!
- AFB: it was a whole saga. But, it was the discrepancy and the interest in it that prompted many very deep QCD developments. Whatever the outcome, we do understand QCD/SM much better now.
- Stop searches (especially stealth stop). Can only be done with high-precision in the SM predictions. Again, this example points at the big picture of possibilities!



Czakon, Mitov, Papucci, Ruderman, Weiler '14 ATLAS '14 (1406.5375)

Few thoughts on BSM and QCD

- Improved gluon pdf (from LHC measurements of NNLO top and dijets) has implications to many processes at the LHC.
 - Example: bSM production at large masses

"Old" (i.e. usual) and "new" (including inclusive NNLO top data) gluon pdf at large x:



... and implied PDF uncertainty due to "old" vs. "new" gluon pdf:

Czakon, Mangano, Mitov, Rojo '13



QCD at LHC: theory developments

Few thoughts on BSM and QCD

- Diboson excess:
 - Understanding better how to search for bumps
 - Tails of distributions (we are searching for bumps on a smooth background; fine but wat is its slope?)
 - Jets and their structure
- Current diphoton excess:
 - It is a clean signal, which is great. But:
 - If it is an extra Higgs, it should decay to tops
 - Why hasn't it been seen in the Mtt spectrum?
 - One needs detailed estimate of effects; likely we are talking about O(5%) effect.
 - With large bins and current errors (even NNLO that just appeared) this is a small effect.
 - One has to devise new strategies for such searches, and this excess (real or not) is an excellent motivation. Work underway.

For QCD aficionados:

NLO automation, parton showers



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NLO calculations: a sample of full(*) automation

Process	Syntax	Cross section (pb)			
Single Higgs product	ion	LO 13 TeV	NLO 13 TeV		
g.1 $pp \rightarrow H$ (HEI g.2 $pp \rightarrow Hj$ (HE	TT) pp>h FT) pp>h j FET) pp>h j		$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$		
$g.3 pp \to \pi j j (\pi)$		$3.020 \pm 0.002 \cdot 10^{-34.7\%} - 1.7\%$	$5.124 \pm 0.020 \cdot 10^{-1.5\%}$		
g.4 $pp \rightarrow Hjj$ (V) g.5 $pp \rightarrow Hjjj$ (V)	BF) pp>hjj\$\$ w+ w-z /BF) pp>hjjj\$\$ w+ w-	$ \begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	$ \begin{array}{cccc} 1.900 \pm 0.006 \cdot 10^{0} & {}^{+ 0.8 \% }_{- 0.9 \% } & {}^{+ 2.0 \% }_{- 0.9 \% } \\ 3.085 \pm 0.010 \cdot 10^{-1} & {}^{+ 2.0 \% }_{- 3.0 \% } & {}^{+ 1.5 \% }_{- 3.0 \% } \end{array} $		
$ \begin{array}{ccc} {\rm g.6} & pp \rightarrow HW^{\pm} \\ {\rm g.7} & pp \rightarrow HW^{\pm} j \\ {\rm g.8^*} & pp \rightarrow HW^{\pm} j \\ \end{array} $	pp>hwpm pp>hwpmj j pp>hwpmjj	$\begin{array}{rrrr} 1.195 \pm 0.002 \cdot 10^{0} & +3.5\% & +1.9\% \\ -4.5\% & -1.5\% \\ 4.018 \pm 0.003 \cdot 10^{-1} & +10.7\% & +1.2\% \\ -9.3\% & -0.9\% \\ 1.198 \pm 0.016 \cdot 10^{-1} & +26.1\% & +0.8\% \\ -19.4\% & -0.6\% \end{array}$	$\begin{array}{rrrr} 1.419 \pm 0.005\cdot10^{0} & +2.1\% & +1.9\% \\ -2.6\% & -1.4\% \\ 4.842 \pm 0.017\cdot10^{-1} & +3.6\% & +1.2\% \\ -3.7\% & -1.0\% \\ 1.574 \pm 0.014\cdot10^{-1} & +5.0\% & +0.9\% \\ -6.5\% & -0.6\% \end{array}$		
$ \begin{array}{ll} {\rm g.9} & pp {\rightarrow} HZ \\ {\rm g.10} & pp {\rightarrow} HZ j \\ {\rm g.11^*} & pp {\rightarrow} HZ jj \end{array} $	p p > h z p p > h z j p p > h z j j	$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	$\begin{array}{rrrr} 7.674 \pm 0.027 \cdot 10^{-1} & +2.0\% & +1.9\% \\ -2.5\% & -1.4\% \\ 2.667 \pm 0.010 \cdot 10^{-1} & +3.5\% & +1.1\% \\ -3.6\% & -0.9\% \\ 8.753 \pm 0.037 \cdot 10^{-2} & +4.8\% & +0.7\% \\ -6.3\% & -0.6\% \end{array}$		
$ \begin{array}{ll} {\rm g.12^*} & pp \rightarrow HW^+W \\ {\rm g.13^*} & pp \rightarrow HW^\pm \gamma \\ {\rm g.14^*} & pp \rightarrow HZW^\pm \\ {\rm g.15^*} & pp \rightarrow HZZ \end{array} $	<pre>r^- (4f) p p > h w+ w- p p > h wpm a p p > h z wpm p p > h z z</pre>	$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$		
$ \begin{array}{ll} {\rm g.16} & pp \rightarrow Ht\bar{t} \\ {\rm g.17} & pp \rightarrow Htj \\ {\rm g.18} & pp \rightarrow Hb\bar{b} \ (4f \end{array} $	p p > h t t~ p p > h tt j p p > h b b~	$ \begin{array}{ccccc} 3.579 \pm 0.003 \cdot 10^{-1} & +30.0\% & +1.7\% \\ -21.5\% & -2.0\% \\ 4.994 \pm 0.005 \cdot 10^{-2} & +2.4\% & +1.2\% \\ -4.2\% & -1.3\% \\ 4.983 \pm 0.002 \cdot 10^{-1} & +28.1\% & +1.5\% \\ -21.0\% & -1.8\% \end{array} $	$\begin{array}{rrrr} 4.608\pm 0.016\cdot 10^{-1} & +5.7\% & +2.0\% \\ -9.0\% & -2.3\% \\ 6.328\pm 0.022\cdot 10^{-2} & +2.9\% & +1.5\% \\ -1.8\% & -1.6\% \\ 6.085\pm 0.026\cdot 10^{-1} & +7.3\% & +1.6\% \\ -9.6\% & -2.0\% \end{array}$		
g.19 $pp \rightarrow Ht\bar{t}j$ g.20* $pp \rightarrow Hb\bar{b}j$ (4	p p > h t t∼ j f) p p > h b b∼ j	$\begin{array}{rrrr} 2.674 \pm 0.041 \cdot 10^{-1} & +45.6\% & +2.6\% \\ & -29.2\% & -2.9\% \\ 7.367 \pm 0.002 \cdot 10^{-2} & +45.6\% & +1.8\% \\ & -29.1\% & -2.1\% \end{array}$	$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$		

MadGraph5_aMC@NLO: sample from 172 processes

Courtesy of M. Grazzini

*) within reason and some limits ...

QCD at LHC: theory developments

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NLO calculations: full(*) automation

- NLO calculations have become so advanced and almost fully automated that, really, there is no excuse to use LO in serious analyses!
- I would mention the aMC@NLO collaboration which has taken the approach of full automation + shower following the extremely successful MC@NLO approach.
- NLO automation allows not only QCD but any SM process. In principle these are contained now in the aMC@NLO.
- Similar developments from the Sherpa+OpenLoops collaboration (see arXiv:1412.5157)
- The number of high-quality works I can't cover here is enormous. Let me only mention few:
 - Denner/Dittmaier et al
 - The Helac collaboration
 - GOSAM project
 - Njet library
 - BlackHat Collaboration
 - MCFM
- Among the most impressive results ever achieved at NLO is the monstrous tt+jet calculation with full off-shell effects and top decay:
 Bevilacqua, Hartanto, Kraus, Worek 1509.09242



QCD at LHC: theory developments

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Summary

Summary

- QCD is the workhorse of LHC physics
- Great recent developments allow for unprecedented accuracy and flexibility:
 - NLO calculations are mature and used everywhere
 - NLO is now fully included in "event generators" like MC@NLO, POWHEG, Sherpa.
 - NNLO is now actively developed and very soon all major 2-to-2 processes (which I discussed here) will be completed.
 - For newest results: stay tuned to the Moriond presentations during next 2 weeks.
- What about the future?
 - Improved accuracies in all interesting processes. Match/beat experimental precision.
 - Ultimately, we want to help answer the question: is there New Physics in the TeV range?
 - People are also thinking about the far future:
 - Future e⁺e⁻ collider of some sort (ILC, CLIC, etc)
 - A future 100 TeV hadron collider (i,.e. much bigger future LHC)