

Cosmological solutions from two-measures model with inflaton field

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The two-measures model Guendelman, Nissimov and Pacheva 2014, 2016

The action of the model $S = S_{darkon} + S_{inflaton}$ is :

$$S_{darkon} = \int d^4x \sqrt{-g} (R(g, \Gamma)) + \int d^4x (\sqrt{-g} + \Phi(C)) L(u, X)$$

$$S_{inflaton} = \int d^4x \Phi_1(A) (R + L^{(1)}) + \int d^4x \Phi_2(B) \left(L^{(2)} + \frac{\Phi(H)}{\sqrt{-g}} \right)$$

where we have auxiliary fields $\Phi_i(X) = \frac{1}{3} \epsilon^{\mu\nu\kappa\lambda} \partial_\mu X_{\nu\kappa\lambda}$ and

$$L(u, X) = -\frac{1}{2} g^{\mu\nu} \partial_\mu u \partial_\nu u - W(u)$$

$$L^{(1)} = -\frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi), \quad V(\phi) = f_1 e^{-\alpha\phi}$$

$$L^{(2)} = -\frac{b}{2} e^{-\alpha\phi} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi + U(\phi), \quad U(\phi) = f_2 e^{-2\alpha\phi}$$

From the equations of motion we obtain the following new constants:

$$L(u, X) = -2M_0 = \text{const}, \quad \frac{\Phi_2(B)}{\sqrt{-g}} = \chi_2 = \text{const}$$

$$R + L^{(1)} = -M_1 = \text{const}, \quad L^{(2)} + \frac{\Phi(H)}{\sqrt{-g}} = -M_2 = \text{const}$$

and the effective potential:

$$U_{\text{eff}}(\phi) = \frac{(V_1(\phi) + M_1)^2}{4\chi_2(U(\phi) + M_2)} \quad \text{with asymptotics} \quad U_- = \frac{f_1^2}{4\chi_2 f_2}, \quad U_+ = \frac{M_1^2}{4\chi_2 M_2}$$

The equations in FLRW metric

The system of equations we need to solve is:

$$v^3 + 3av + 2b = 0 \text{ for } a = \frac{-1}{3} \frac{V(\phi) + M_1 - \frac{1}{2}\chi_2 b e^{-\alpha\phi} \dot{\phi}^2}{\chi_2(U(\phi) + M_2) - 2M_0}, b = \frac{-p_u}{2a(t)^3(\chi_2(U(\phi) + M_2) - 2M_0)} \quad (1)$$

$$\dot{a}(t) = \sqrt{\frac{\rho}{6}}a(t), \quad \rho = \frac{1}{2}\dot{\phi}^2\left(1 + \frac{3}{4}\chi_2 b e^{-\alpha\phi} v^2\right) + \frac{v^2}{4}(V + M_1) + \frac{3p_u v}{4a(t)^3} \quad (2)$$

$$\frac{d}{dt} \left(a(t)^3 \dot{\phi} \left(1 + \frac{\chi_2}{2} b e^{-\alpha\phi} v^2\right) \right) + a(t)^3 \left(\alpha \frac{\dot{\phi}^2}{4} \chi_2 b e^{-\alpha\phi} v^2 + \frac{1}{2} V_\phi v^2 - \chi_2 U_\phi \frac{v^4}{4} \right) = 0 \quad (3)$$

The parameters of this system are 12:

$$\{\alpha, b_0, M_0, M_1, M_2, f_1, f_2, p_u, \chi_2\}$$

Initial and boundary conditions:

I. $a(0) = 10^{-12}, \phi(0) = \phi_0, \dot{\phi}(0) = 0$ – initial conditions

II. $a(1) = 1$ – normalization

III. $a''(t) = 0$ – in 3 points

Constraints on the parameters:

$$f_1^2/f_2 \gg M_1^2/M_2$$

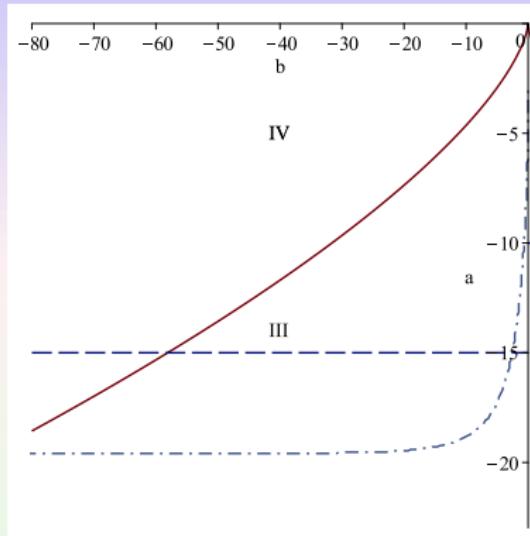
In Planck units: $M_{Pl} = \sqrt{2}$:

$$M_1 = 4 \cdot 10^{-60}, M_2 = 4, f_1 \sim 10^{-8}, f_2 \sim 10^{-8}$$

The evolution of the Universe in the $[a, b]$ plane

Staicova & Stoilov, Mod. Phys. Lett. A, 32, 1 (2017)

Staicova & Stoilov, arXiv:1801.07133



Evolution of the parameters $[a(t), b(t)]$ for the darkon (dash) and the inflaton (dot-dash).

The evolution of the Universe starts from $b \rightarrow -\infty$ and finishes at $b \rightarrow 0$

We have chosen the parameters in such a way that: $b = \frac{-\rho_u}{2a(t)^3(\chi_2(U+M_2)-2M_0)} < 0$.

Numerical solutions

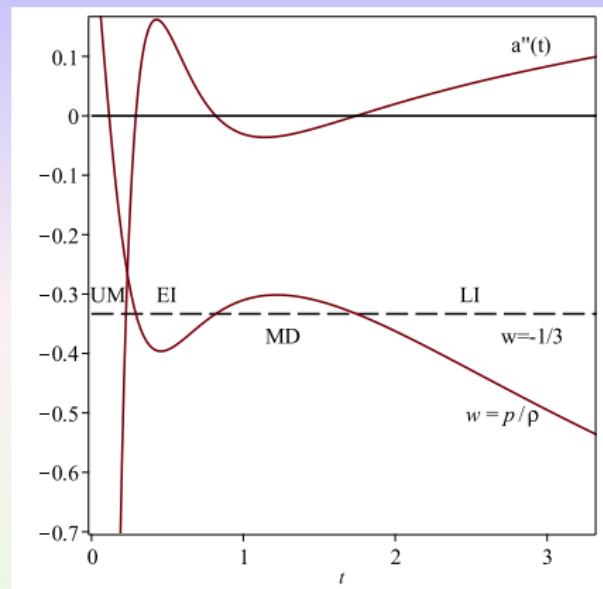
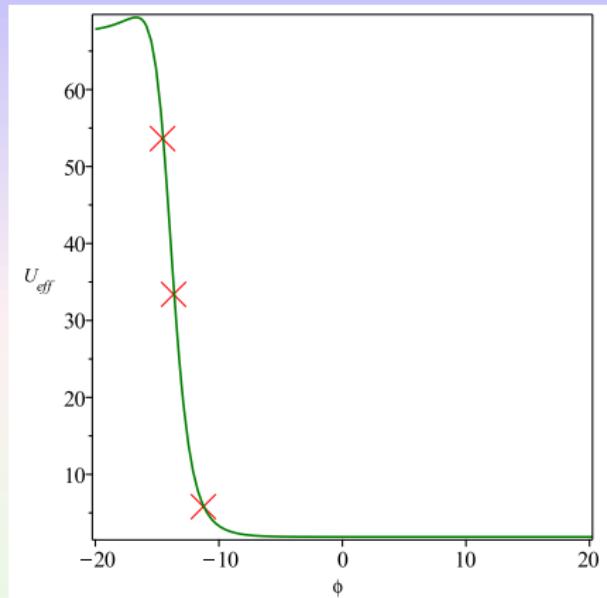
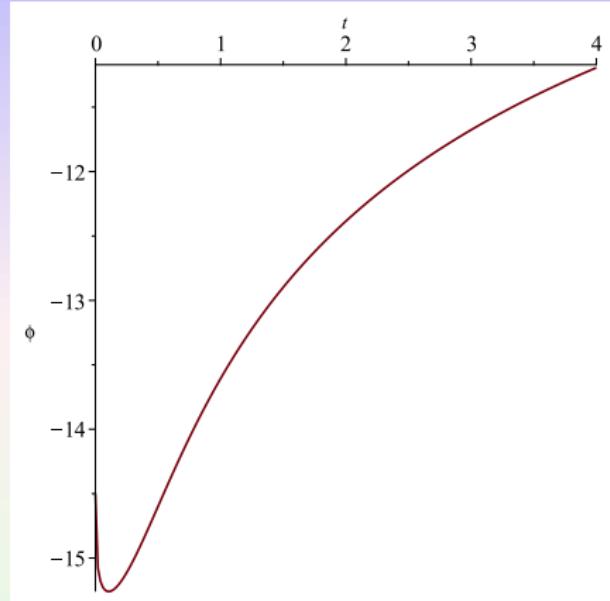
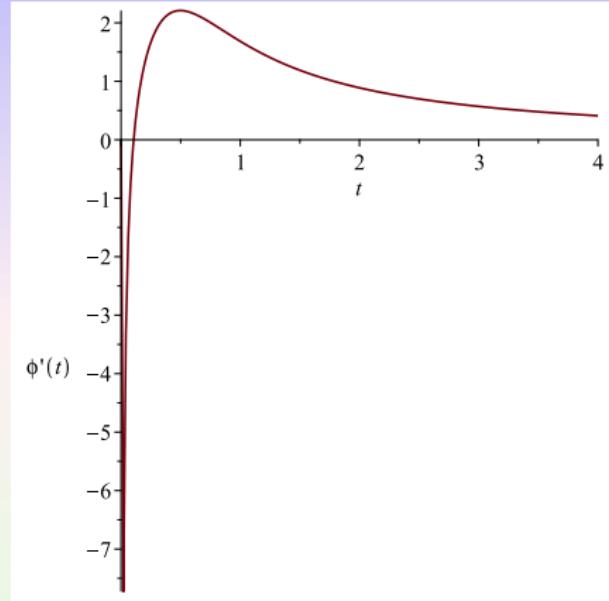


Figure: a) The effective potential.
b) $\ddot{a}(t)$ and $w = p/\rho$, where UM is ultra-relativistic matter domination , EI – the early inflation, MD – the matter domination (MD) and LI – the late inflation.

Strong friction term on ϕ !



(a) $\phi(t)$



(b) $\dot{\phi}(t)$

From asymptotical analysis: assuming $a(t) = e^{Ht}$, we obtain:

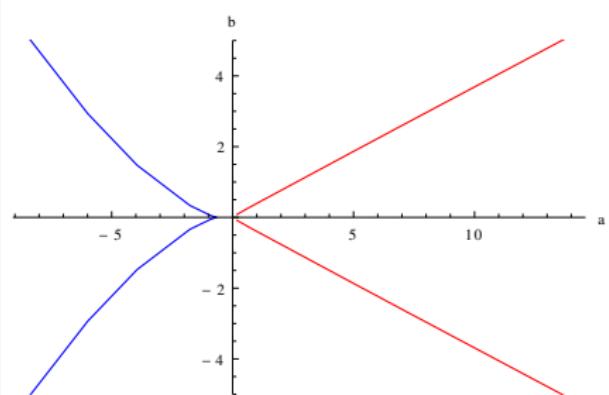
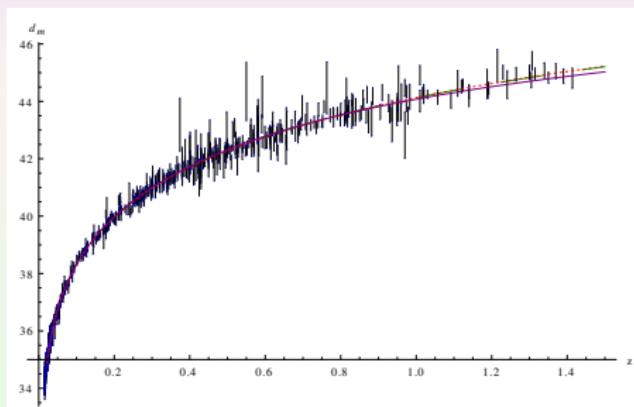
$$\phi(t) = C_1 + C_2 e^{-3Ht}$$

Conclusions:

1. We could not confirm the parameter estimation in the original work.
2. It is not possible to start from the left plateau and to obtain physically realistic solutions
3. Friction term which stops the evolution of the inflaton
4. The theory can produce “realistic” Universe if evolution starts from the slope.

Staicova & Stoilov, Mod. Phys. Lett. A, **32**, 1 (2017)

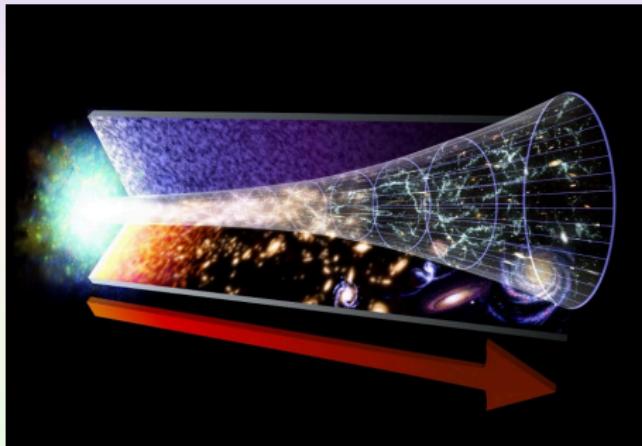
Staicova & Stoilov, arXiv:1801.07133



Fitting the darkon only model with data from Supernovae Type Ia.

Thank you for your attention!

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