## The quantum Fourier transform and its applications

to successive powers of two. The final state of the first register is easily seen to be:

$$\frac{1}{2^{t/2}} \left( |0\rangle + e^{2\pi i 2^{t-1} \varphi} |1\rangle \right) \left( |0\rangle + e^{2\pi i 2^{t-2} \varphi} |1\rangle \right) \dots \left( |0\rangle + e^{2\pi i 2^{0} \varphi} |1\rangle \right) 
= \frac{1}{2^{t/2}} \sum_{k=0}^{2^{t-1}} e^{2\pi i \varphi k} |k\rangle.$$
(5.20)

We omit the second register from this description, since it stays in the state  $|u\rangle$  throughout the computation.

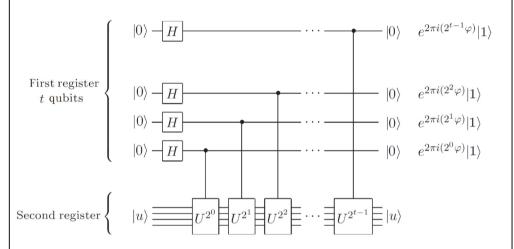


Figure 5.2. The first stage of the phase estimation procedure. Normalization factors of  $1/\sqrt{2}$  have been omitted, on the right.

$$\frac{1}{2^{t/2}} \left( |0\rangle + e^{2\pi i 0.\varphi_t} |1\rangle \right) \left( |0\rangle + e^{2\pi i 0.\varphi_{t-1}\varphi_t} |1\rangle \right) \dots \left( |0\rangle + e^{2\pi i 0.\varphi_1 \varphi_2 \dots \varphi_t} |1\rangle \right) . (5.21)$$

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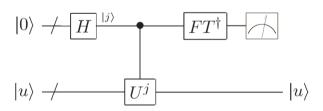
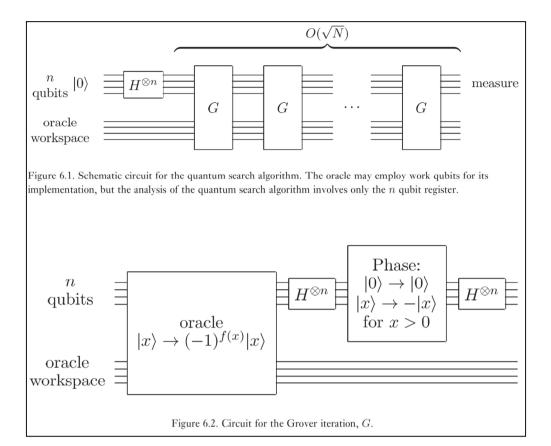


Figure 5.3. Schematic of the overall phase estimation procedure. The top t qubits (the '/' denotes a bundle of wires, as usual) are the first register, and the bottom qubits are the second register, numbering as many as required to perform U.  $|u\rangle$  is an eigenstate of U with eigenvalue  $e^{2\pi i \varphi}$ . The output of the measurement is an approximation to  $\varphi$  accurate to  $t - \left\lceil \log\left(2 + \frac{1}{2\epsilon}\right) \right\rceil$  bits, with probability of success at least  $1 - \epsilon$ .

http://theo.inrne.bas.bg/~mitov/qi21/notes\_df34g2tb3/NC2010.pdf



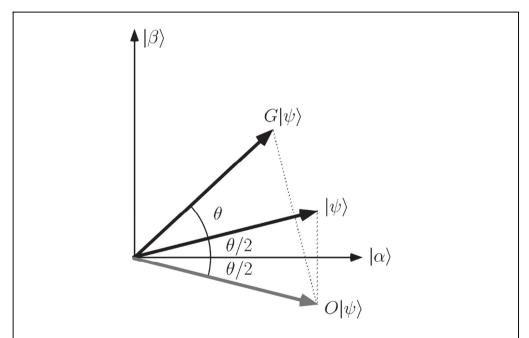


Figure 6.3. The action of a single Grover iteration, G: the state vector is rotated by  $\theta$  towards the superposition  $|\beta\rangle$  of all solutions to the search problem. Initially, it is inclined at angle  $\theta/2$  from  $|\alpha\rangle$ , a state orthogonal to  $|\beta\rangle$ . An oracle operation O reflects the state about the state  $|\alpha\rangle$ , then the operation  $2|\psi\rangle\langle\psi|-I$  reflects it about  $|\psi\rangle$ . In the figure  $|\alpha\rangle$  and  $|\beta\rangle$  are lengthened slightly to reduce clutter (all states should be unit vectors). After repeated Grover iterations, the state vector gets close to  $|\beta\rangle$ , at which point an observation in the computational basis outputs a solution to the search problem with high probability. The remarkable efficiency of the algorithm occurs because  $\theta$  behaves like  $\Omega(\sqrt{M/N})$ , so only  $O(\sqrt{N/M})$  applications of G are required to rotate the state vector close to  $|\beta\rangle$ .