New Integrable RG Flows with Parafermions

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Based on

Z. Bajnok & CA [arXiv:2407.06582]

Integrability, Q-systems and Cluster Algebras, Varna, 2024

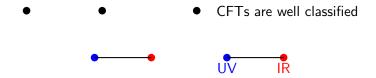
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CFTs are well classified

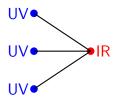




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Can we classify all possible 2D QFTs which flow into a given IR CFT?

Plan of Talk

- 1. New Approach [LeClair & CA (2022)]
 - Classify UV CFTs connected to an IR CFT
 - But only UV CFTs can be identified
- 2. Identifying QFTs [Bajnok & CA (2024)]
 - Focus on the unitary minimal CFTs \mathcal{M}_p
 - Find and identify all QFTs which flow into this
- 3. Concluding remarks

PART 1. RG flows from IR to UV

Conventional Approach (top down)

- UV CFT+a relevant field

 IR CFT+ irrelevant fields
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 - Common in CMP since physics at low T is important to find non-trivial (Wilson-Fisher) fixed points

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- Many more RG flows have been conjectured
 - have been guessed based on conjectured TBA or NLIE
 - Lagrangians / exact S-matrices are missing



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- Use $T\overline{T}$ as a ladder

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- Preserve integrability
- Exact results possible for even non-integrable theories

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New IQFTs = an IQFT +
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 \bullet Exact S-matrices are given by additional CDD factors

$$\mathsf{S}(\theta) = \prod_{s=1}^{\infty} e^{2i\alpha_s M^{2s} \sinh(s\theta)} \cdot \mathsf{S}_0(\theta)$$

• Burgers Equation $(s = 1) [CFT + \alpha_1[T\overline{T}]_1]$

$$\partial_{lpha} E + E \partial_R E = 0 \quad o \quad c_{ ext{eff}}(R) = rac{2c_0}{1 + \sqrt{1 + rac{2\pilpha_1c_0}{3R^2}}}$$

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- We show that the singularities can be avoided if we fine-tune all α_s

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S-matrices of CFTs from a massless limit

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- S-matrices between RL, LR are trivial since $\theta_{12} \to \pm \infty$

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Summary

$$\boxed{\mathsf{S}^{RR}(\theta) = \mathsf{S}^{LL}(\theta) = \mathsf{S}_0, \; \mathsf{S}^{RL}(\theta) = \mathsf{S}^{LR}(-\theta) = \mathsf{S}_{\mathsf{CDD}}(\theta)}$$

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from which S_{CDD} are arbitrary products of only two factors

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- We can work with UV limit of TBA, namely, "plateaux equations"

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• $c_{\rm eff}$ in terms of some dilog identities

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• From the Plateaux equations, we find only

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- We analyzed other group G = su(3), su(4) etc to classify all possible UV CFTs based on central charges
- But central charges are not enough to identify the QFTs

PART 2. Identifying UV complete QFTs

with Z. Bajnok

• Consider k = p - 2 with G = su(2)

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Particle spectrum: massive kinks

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• S-matrix of kinks $S_{\rm RSOS}^{\rm [p]}(\theta)$: non-diagonal [Bernard, LeClair (1990)]

$$K_{ab}(\theta_1)K_{bc}(\theta_2) \quad \rightarrow \quad K_{ad}(\theta_2)K_{dc}(\theta_1)$$

$$\mathsf{S}_p(\theta)_{dc}^{ab} = \mathit{U}(\theta) \; (X_{db}^{ac})^{\frac{i\theta}{2\pi}} \left[(X_{db}^{ac})^{\frac{1}{2}} \sinh \left(\frac{\theta}{p} \right) \; \delta_{db} + \sinh \left(\frac{i\pi - \theta}{p} \right) \; \delta_{ac} \right]$$

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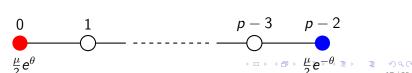
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Derived the conjectured TBA (partially) shown before

$$\epsilon_{\mathsf{a}}(heta) = rac{\mu R}{2} (\delta_{\mathsf{a}0} e^{ heta} + \delta_{\mathsf{a},p-2} e^{- heta}) - \sum_{b=0}^{p-2} \mathbb{I}_{\mathsf{a}b} arphi \star \log[1 + e^{-\epsilon_b}](heta)$$



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• $[T\overline{T}]_s$ deformations

$$\mathcal{M}_p + \lambda' \, \Phi_{3,1} + \sum_{s \geq 1} \alpha_s \, [\mathcal{T} \, \overline{\mathcal{T}}]_s$$

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Introduce CDD factors to RL, LR sectors

$$\tilde{\mathsf{S}}^{RL}_{p}(\theta) = \mathsf{S}_{\mathsf{CDD}} \cdot \mathsf{S}^{RL}_{p}(\theta)$$

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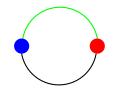
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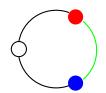
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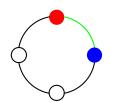
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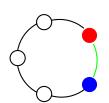
$$\tilde{\mathsf{S}}_p^{RL}(\theta) = \mathsf{S}_{\mathsf{CDD}} \cdot \mathsf{S}_p^{RL}(\theta)$$

• Being diagonal, S_{CDD} introduces additional kernels between $R\ \&\ L$









Plateaux equations

$$x_n = (1+x_{n-1})^{1/2}(1+x_{n+1})^{1/2}, \quad n=1,\cdots,p-3,$$

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• UV complete only with $\tilde{k} = \frac{1}{2}$ which means

$$\mathsf{S}_{\mathsf{CDD}}^{(1)} = - anh \left(rac{ heta - eta}{2} - rac{i\pi}{4}
ight) \quad o \quad arphi_{\mathit{RL}} = rac{1}{\cosh(heta - eta)}$$

• Plateaux equations

$$x_n = (1+x_{n-1})^{1/2}(1+x_{n+1})^{1/2}, \quad n=1,\cdots,p-3,$$

 $x_0 = (1+x_1)^{1/2}(1+x_{p-2})^{\tilde{k}}, \quad x_{p-2} = (1+x_{p-3})^{1/2}(1+x_0)^{\tilde{k}}$

• UV complete only with $\tilde{k} = \frac{1}{2}$ which means

$$S_{CDD}^{(1)} = -\tanh\left(\frac{\theta - \beta}{2} - \frac{i\pi}{4}\right) \quad \rightarrow \quad \varphi_{RL} = \frac{1}{\cosh(\theta - \beta)}$$

• We also notice that only real β (p > 3) becomes UV complete and

$$c_{\text{eff}} = 3 \frac{p-1}{p+1}$$

• sinh-Gordon model with ${\it N}=1$ super [Kim, Rim, Zamolodchikov, CA (2002)]

$$\mathcal{L} = \mathrm{Kin.} - rac{i}{2} \psi \overline{\psi} W''(\phi) + 2\pi [W'(\phi)]^2, \quad W(\phi) = -\mu \sinh(b\phi)$$

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- TBA is given by

$$\epsilon^{R}(\theta) = \frac{\mu R}{2} e^{\theta} - \varphi^{RL} \star \log[1 + e^{-\epsilon^{L}}],$$

$$\varphi^{RL}(\theta) = \frac{1}{\cosh(\theta - ia)} + \frac{1}{\cosh(\theta + ia)}, \quad a = \frac{1 - b^{2}}{1 + b^{2}}$$



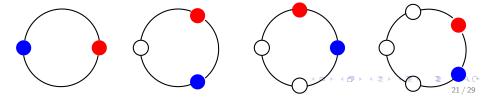
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PshG model

$$\mathcal{L}_{\mathsf{PShG}} = \mathcal{L}_{\mathsf{PF}} + rac{1}{4\pi} (\partial_{\mu}\phi)^2 - \kappa \left(\psi_1 \overline{\psi}_1 \mathsf{e}^{2b\phi} + \eta \, \psi_1^{\dagger} \overline{\psi}_1^{\dagger} \mathsf{e}^{-2b\phi}
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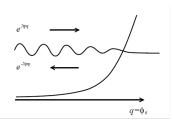
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- These types of QFTs compute $c_{\rm eff}$ from momentum quantization using Reflection amplitudes

• Primary fields $e^{2\alpha\phi}$ and $e^{2(Q-\alpha)\phi}$ identified upto constant

$$e^{2(\frac{Q}{2}+ip)\phi}=\mathcal{R}(p)e^{2(\frac{Q}{2}-ip)\phi}, \qquad \alpha=\frac{Q}{2}+ip, \ Q=b+\frac{1}{kb}$$

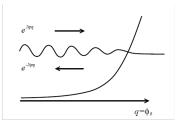
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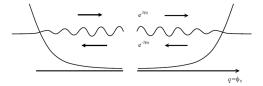
 This amplitude has been computed in the PLFT [Baseilhac, Fateev (1998)]

Momentum quantization condition

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the perturbation introduces another wall

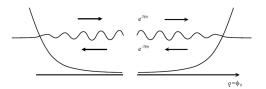
$$\delta^{(k)}(p) = \pi + 4Qp \ln \frac{R}{2\pi} = \sum_{n=\mathrm{odd}} \delta_n p^n, \quad \mathcal{R}(p) = e^{i\delta^{(k)}(p)}$$



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Scaling function can be obtained from

$$c_{\text{eff}} = \frac{3k}{k+2} - 24p^2 + \mathcal{O}(R)$$

$$= \frac{3k}{k+2} - \frac{3\pi^2}{2Q^2 \ln^2 x} - \frac{3\pi^2 \delta_1}{4Q^3 \ln^3 x} - \frac{9\pi^2 \delta_1^2}{32Q^4 \ln^4 x} - \frac{3(2\pi^2 \delta_1^3 + \pi^4 \delta_3)}{64Q^5 \ln^5 x} + \dots \quad (x = \frac{R}{2\pi})$$

Reflection amplitude vs. massless TBA

Need to solve numerically with high precision

$$c_{\text{eff}}(r) = \frac{3k}{k+2} + \sum_{n=2} \frac{c_n(k)}{(\log r)^n} + O(r)$$

• Two match very well at self-dual coupling $b=\frac{1}{\sqrt{k}}$

	<i>c</i> ₂	<i>c</i> ₃	C4	<i>C</i> ₅
k=2	7.402199	42.7620	185.218	714.563
	7.402203	42.7628	185.282	717.247
k = 3	11.10332	77.8573	409.598	1924.84
	11.10330	77.8543	409.425	1919.36
k = 4	14.8045	119.428	722.797	3898.75
	14.8044	119.4197	722.475	3892.55

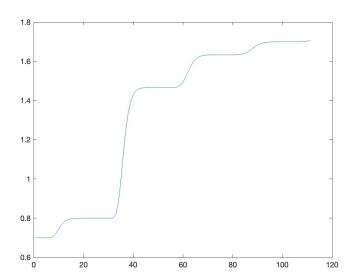
What if $\beta \neq 0$ for generic $p \geq 3$?

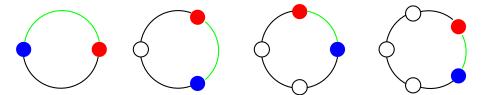
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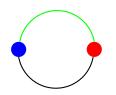
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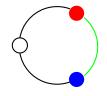
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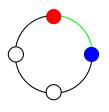
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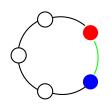










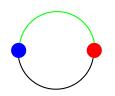


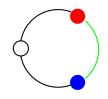
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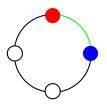
$$\mathcal{M}_{k+1} \to \mathbb{Z}_k \mathcal{M}_1 \to \mathbb{Z}_k \mathcal{M}_{k+1} \to \mathbb{Z}_k \mathcal{M}_{2k+1}$$

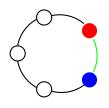
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Reproduce PF staircase TBA conjectured by [Dorey, Ravanini (1993)]

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- Connection with non-invertible symmetries

Thank you for attention!